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DOI: [https://doi.org/10.1007/JHEP11\(2010\)102](https://doi.org/10.1007/JHEP11(2010)102)

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ZORA URL: <https://doi.org/10.5167/uzh-41359>

Journal Article

Accepted Version

Originally published at:

Gehrmann, T; Glover, E W N; Huber, T; Ikizlerli, N; Studerus, C (2010). The quark and gluon form factors to three loops in QCD through to $\mathcal{O}(\epsilon^2)$. *Journal of High Energy Physics*, (11):012.

DOI: [https://doi.org/10.1007/JHEP11\(2010\)102](https://doi.org/10.1007/JHEP11(2010)102)

The quark and gluon form factors to three loops in QCD through to $\mathcal{O}(\epsilon^2)$

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ABSTRACT: We give explicit formulae for the $\mathcal{O}(\epsilon)$ and $\mathcal{O}(\epsilon^2)$ contributions to the unrenormalised three loop QCD corrections to quark and gluon form factors. These contributions have at most transcendentality weight eight. The $\mathcal{O}(\epsilon)$ terms of the three-loop form factors are required for the extraction of the four-loop quark and gluon collinear anomalous dimensions. The $\mathcal{O}(\epsilon^2)$ terms represent an irreducible contribution to the finite part of the form factors at four-loops. For the sake of completeness, we also give the contributions to the one and two loop form factors to the same transcendentality weight eight.

KEYWORDS: QCD, Multi-loop calculations.

The form factors are fundamental ingredients for many precision calculations in QCD. These basic building blocks describe the coupling of an external, colour-neutral off-shell particle to a pair of partons: the quark form factor is the coupling of a virtual photon to a quark-antiquark pair, while the gluon form factor is the coupling of a Higgs boson to a pair of gluons through an effective Lagrangian.

The form factors are phenomenologically important and appear directly as virtual higher-order corrections in coefficient functions for the inclusive Drell-Yan process [1–3] and the inclusive Higgs production cross section [3–6]. The form factors also display a non-trivial infrared pole structure which is determined by the infrared factorisation formula. This implies that their infrared pole coefficients can be used to extract fundamental constants such as the cusp anomalous dimensions which control the structure of soft divergences and the collinear quark and gluon anomalous dimensions. In fact, the cusp anomalous dimensions were first obtained to three loops from the asymptotic behaviour of splitting functions [7, 8]. However, it was the calculation [9, 10] of the pole terms of the three-loop form factors (and finite plus subleading terms in the two-loop and one-loop form factors [11–13]), which led to the derivation of the three-loop collinear anomalous dimensions [9, 14, 15].

The infrared factorisation formula for a given form factor (or more generally for a given multi-leg amplitude) at a certain number of loops involves infrared singularity operators acting on the form factor evaluated with a lower number of loops. These infrared singularity operators contain explicit infrared poles $1/\epsilon^2$ and $1/\epsilon$. Therefore, the computation of the finite contribution to any n -loop form factor relies on contributions from $(n - m)$ -loops evaluated to $\mathcal{O}(\epsilon^{2m})$.

At present, the state of the art is at the three-loop level for the massless quark and gluon form factors. There are 22 master integrals shown in Fig. 1, of which 14 are genuine three-loop vertex functions ($A_{t,i}$ -type), 4 are three-loop propagator integrals ($B_{t,i}$ -type) and 4 are products of one-loop and two-loop integrals ($C_{t,i}$ -type). In this notation, the index t denotes the number of propagators, and i is simply enumerating the topologically different integrals with the same number of propagators. Expressions for the form factors in terms of the 22 independent master integrals, and valid for any value of the dimension D , are given in Ref. [16]. The $B_{t,i}$ -type integrals were computed to finite order in [17, 18] and supplemented by the higher order terms in [19]. Explicit expansions of the $A_{t,i}$ -type integrals were obtained in Refs. [20–23] using Mellin-Barnes techniques. They enabled the evaluation of the three-loop form factors up to and including the finite contributions [16, 23, 24]. The deepest pole contribution is of $\mathcal{O}(1/\epsilon^6)$. Correspondingly, the finite terms are of at most transcendentality weight six, that is terms such as π^6 (ζ_2^3) or ζ_3^2 .

More recently [25], 20 of the three-loop master integrals have been re-evaluated up to transcendentality weight eight using dimensional recurrence relations [26] and analytic properties of Feynman integrals (the DRA method [27]). Expressions for the two remaining integrals, $B_{8,1}$ and $C_{8,1}$, can be obtained from Refs. [27] and [13] respectively. Once the same normalisation and basis set of multiple zeta values is used, Ref. [25] confirms the earlier result of Ref. [28] for $A_{6,2}$. On the other hand, we confirm a certain subset of master integrals ($B_{6,2}$, $B_{8,1}$, $A_{7,3}$, $A_{7,5}$, $A_{8,1}$, $A_{9,1}$, $A_{9,2}$, $A_{9,4}$) from [25, 27] up to coefficients

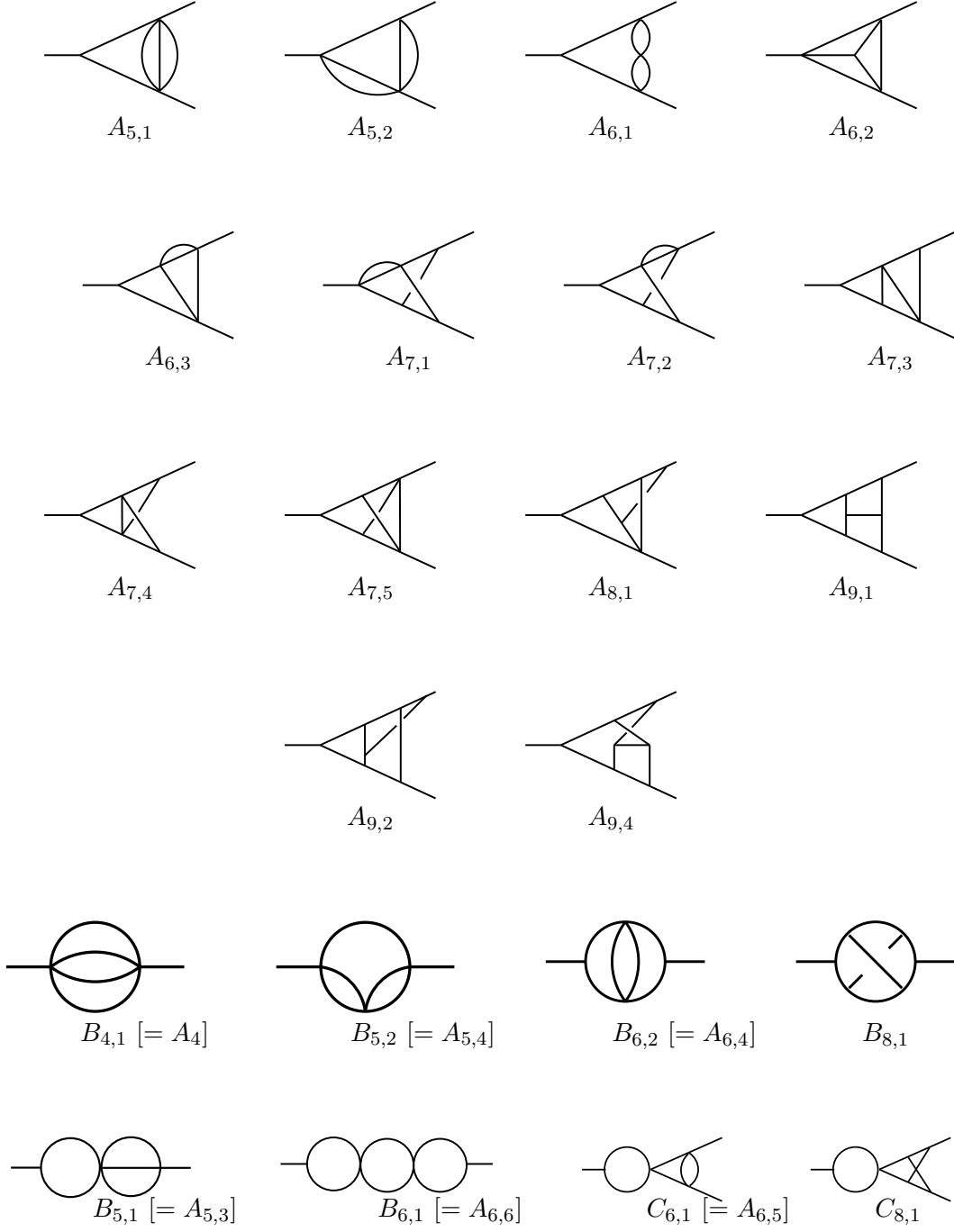


Figure 1: Master integrals for the three-loop form factors. Labels in brackets indicate the naming convention of Ref. [25].

corresponding to weight eight numerically to a precision of one per-mille or better using **MB.m** [29] and **FIESTA** [30, 31]. All other of the 22 master integrals we even confirm analytically through to weight eight by expanding the closed form in terms of hypergeometric functions given in [20, 21] using the **HypExp** package [32].

All master integrals are therefore known up to transcendentality weight eight i.e. terms including π^8 (ζ_2^4), $\zeta_2\zeta_3^2$, $\zeta_3\zeta_5$ as well as the multiple zeta value $\zeta_{5,3}$ (or equivalently $\zeta_{-6,-2}$). The multiple zeta values are defined by (see e.g. [33] and references therein)

$$\zeta(m_1, \dots, m_k) = \sum_{i_1=1}^{\infty} \sum_{i_2=1}^{i_1-1} \cdots \sum_{i_k=1}^{i_{k-1}-1} \prod_{j=1}^k \frac{\text{sgn}(m_j)^{i_j}}{i_j^{|m_j|}}. \quad (1)$$

Specifically, $\zeta_{-6,-2}$ is related to $\zeta_{5,3}$ by [33, 34]

$$\zeta_{-6,-2} = \frac{9}{20}\zeta_{5,3} - \frac{3}{2}\zeta_5\zeta_3 + \frac{781}{4032000}\pi^8. \quad (2)$$

The numerical values of the transcendental constants up to weight eight are:

$$\begin{aligned} \zeta_3 &= 1.2020569031595942854\dots, & \zeta_5 &= 1.0369277551433699263\dots, \\ \zeta_7 &= 1.0083492773819228268\dots, & \zeta_{5,3} &= 0.037707672984847544011\dots \end{aligned}$$

The new results on the higher order terms in the master integrals enable the computation of the three-loop form factors through to $\mathcal{O}(\epsilon^2)$ which is an intrinsic component for the four-loop evaluation of the form factors. This is the topic of this Letter and we give explicit formulae for the $\mathcal{O}(\epsilon)$ and $\mathcal{O}(\epsilon^2)$ contributions to the unrenormalised three loop form factors.

The form factors are the basic vertex functions of an external off-shell current (with virtuality $q^2 = s_{12}$) coupling to a pair of partons with on-shell momenta p_1 and p_2 . One distinguishes time-like ($s_{12} > 0$, i.e. with partons both either in the initial or in the final state) and space-like ($s_{12} < 0$, i.e. with one parton in the initial and one in the final state) configurations. The form factors are described in terms of scalar functions by contracting the respective vertex functions (evaluated in dimensional regularization with $D = 4 - 2\epsilon$ dimensions) with projectors. For massless partons, the full vertex function is described with only a single form factor.

The quark form factor is obtained from the photon-quark-antiquark vertex $\Gamma_{q\bar{q}}^\mu$ by

$$\mathcal{F}^q = -\frac{1}{4(1-\epsilon)q^2} \text{Tr} (p_2 \Gamma_{q\bar{q}}^\mu p_1 \gamma_\mu), \quad (3)$$

while the gluon form factor relates to the effective Higgs-gluon-gluon vertex $\Gamma_{gg}^{\mu\nu}$ as

$$\mathcal{F}^g = \frac{p_1 \cdot p_2 g_{\mu\nu} - p_{1,\mu} p_{2,\nu} - p_{1,\nu} p_{2,\mu}}{2(1-\epsilon)} \Gamma_{gg}^{\mu\nu}. \quad (4)$$

The form factors are expanded in perturbative QCD in powers of the coupling constant, with each power corresponding to a virtual loop. We denote the unrenormalized form factors by \mathcal{F}^a and the renormalized form factors by F^a with $a = q, g$.

At tree level, the Higgs boson does not couple either to the gluon or to massless quarks. In higher orders in perturbation theory, heavy quark loops introduce a coupling between the Higgs boson and gluons. In the limit of infinitely massive quarks, these loops give rise

to an effective Lagrangian [35] mediating the coupling between the scalar Higgs field and the gluon field strength tensor:

$$\mathcal{L}_{\text{int}} = -\frac{\lambda}{4} H F_a^{\mu\nu} F_{a,\mu\nu} . \quad (5)$$

The coupling λ has inverse mass dimension. It can be computed by matching [36, 37] the effective theory to the full standard model cross sections [5].

Direct evaluation of the Feynman diagrams at the appropriate loop order yields the bare (unrenormalised) form factors,

$$\mathcal{F}_b^q(\alpha_s^b, s_{12}) = 1 + \sum_{n=1}^{\infty} \left(\frac{\alpha_s^b}{4\pi} \right)^n \left(\frac{-s_{12}}{\mu_0^2} \right)^{-n\epsilon} S_\epsilon^n \mathcal{F}_n^q, \quad (6)$$

$$\mathcal{F}_b^g(\alpha_s^b, s_{12}) = \lambda^b \left(1 + \sum_{n=1}^{\infty} \left(\frac{\alpha_s^b}{4\pi} \right)^n \left(\frac{-s_{12}}{\mu_0^2} \right)^{-n\epsilon} S_\epsilon^n \mathcal{F}_n^g \right), \quad (7)$$

where μ_0^2 is the mass parameter introduced in dimensional regularisation to maintain a dimensionless coupling in the bare Lagrangian density and where

$$S_\epsilon = e^{-\epsilon\gamma}(4\pi)^\epsilon, \quad \text{with the Euler constant } \gamma = 0.5772\dots \quad (8)$$

The one-loop and two-loop form factors were computed in many places in the literature [9–13]. All-order expressions in terms of one-loop and two-loop master integrals are given in [13]. Explicit expressions for the one- and two-loop form factors through to $\mathcal{O}(\epsilon^5)$ and $\mathcal{O}(\epsilon^3)$ respectively are given already in [16]. To determine the finite piece at the four-loop level, these form factors are needed to one higher power in ϵ , and for the sake of completeness, we quote them here. At one-loop,

$$\begin{aligned} \mathcal{F}_1^q &= \mathcal{F}_1^q|_{\frac{1}{\epsilon^2}} + \dots + \mathcal{F}_1^q|_{\epsilon^5} \\ &+ C_F \left[\epsilon^6 \left(-512 + \frac{381\zeta_7}{7} + \frac{496\zeta_5}{5} + \frac{448\zeta_3}{3} - \frac{434\zeta_3\zeta_5}{15} - \frac{196\zeta_3^2}{9} + 64\zeta_2 \right. \right. \\ &\quad \left. \left. - \frac{93\zeta_2\zeta_5}{10} - \frac{56\zeta_2\zeta_3}{3} + \frac{49\zeta_2\zeta_3^2}{18} + \frac{188\zeta_2^2}{5} - \frac{329\zeta_2^2\zeta_3}{40} + \frac{949\zeta_2^3}{70} + \frac{55779\zeta_2^4}{11200} \right) \right] \\ \mathcal{F}_1^g &= \mathcal{F}_1^g|_{\frac{1}{\epsilon^2}} + \dots + \mathcal{F}_1^g|_{\epsilon^5} \\ &+ C_A \left[\epsilon^6 \left(-126 + \frac{62\zeta_5}{5} + \frac{98\zeta_3}{3} - \frac{434\zeta_3\zeta_5}{15} + 15\zeta_2 - \frac{7\zeta_2\zeta_3}{3} + \frac{49\zeta_2\zeta_3^2}{18} \right. \right. \\ &\quad \left. \left. + \frac{141\zeta_2^2}{20} + \frac{55779\zeta_2^4}{11200} \right) \right] \end{aligned} \quad (9)$$

and at two-loops

$$\begin{aligned} \mathcal{F}_2^q &= \mathcal{F}_2^q|_{\frac{1}{\epsilon^4}} + \dots + \mathcal{F}_2^q|_{\epsilon^3} \\ &+ C_F^2 \left[\epsilon^4 \left(+\frac{637631}{128} - 528\zeta_{5,3} + \frac{27204\zeta_7}{7} - \frac{34001\zeta_5}{10} - \frac{481913\zeta_3}{24} + \frac{33248\zeta_3\zeta_5}{15} \right. \right. \end{aligned}$$

$$\begin{aligned}
& + \frac{36359\zeta_3^2}{9} + \frac{95559\zeta_2}{32} - 198\zeta_2\zeta_5 + \frac{2257\zeta_2\zeta_3}{2} - \frac{4576\zeta_2\zeta_3^2}{9} - \frac{248023\zeta_2^2}{80} \\
& + \frac{5109\zeta_2^2\zeta_3}{5} + \frac{55623\zeta_2^3}{140} + \frac{653901\zeta_2^4}{700} \Bigg) \\
& + C_F C_A \left[\epsilon^4 \left(-\frac{11630115085}{839808} + 264\zeta_{5,3} - \frac{11980\zeta_7}{21} + \frac{1214029\zeta_5}{270} + \frac{84520897\zeta_3}{5832} \right. \right. \\
& \quad - \frac{8266\zeta_3\zeta_5}{5} - \frac{229042\zeta_3^2}{81} - \frac{58499773\zeta_2}{23328} - \frac{829\zeta_2\zeta_5}{15} - \frac{94931\zeta_2\zeta_3}{162} + \frac{3029\zeta_2\zeta_3^2}{9} \\
& \quad \left. \left. + \frac{14915741\zeta_2^2}{6480} - \frac{66379\zeta_2^2\zeta_3}{90} + \frac{4843\zeta_2^3}{30} - \frac{75242\zeta_2^4}{175} \right) \right] \\
& + C_F N_F \left[\epsilon^4 \left(+\frac{996726245}{419904} - \frac{2186\zeta_7}{21} - \frac{42713\zeta_5}{135} - \frac{1951625\zeta_3}{2916} + \frac{4732\zeta_3^2}{81} + \frac{2877653\zeta_2}{11664} \right. \right. \\
& \quad \left. \left. - \frac{242\zeta_2\zeta_5}{15} - \frac{4589\zeta_2\zeta_3}{81} - \frac{309181\zeta_2^2}{3240} + \frac{533\zeta_2^2\zeta_3}{45} - \frac{127\zeta_2^3}{3} \right) \right] \\
\mathcal{F}_2^g &= \mathcal{F}_2^g|_{\frac{1}{\epsilon^4}} + \dots + \mathcal{F}_2^g|_{\epsilon^3} \\
& + C_A^2 \left[\epsilon^4 \left(+\frac{1371828689}{209952} - 264\zeta_{5,3} + \frac{56155\zeta_7}{42} - \frac{161266\zeta_5}{135} - \frac{5108944\zeta_3}{729} + \frac{1690\zeta_3\zeta_5}{3} \right. \right. \\
& \quad + \frac{85559\zeta_3^2}{81} - \frac{219275\zeta_2}{1944} - \frac{1001\zeta_2\zeta_5}{5} + \frac{11858\zeta_2\zeta_3}{27} - \frac{1547\zeta_2\zeta_3^2}{9} - \frac{187733\zeta_2^2}{180} \\
& \quad \left. \left. + \frac{22781\zeta_2^2\zeta_3}{90} + \frac{123079\zeta_2^3}{1260} + \frac{50419\zeta_2^4}{100} \right) \right] \\
& + C_A N_F \left[\epsilon^4 \left(-\frac{232282297}{104976} + \frac{229\zeta_7}{21} - \frac{24518\zeta_5}{135} - \frac{301886\zeta_3}{729} + \frac{22060\zeta_3^2}{81} \right. \right. \\
& \quad \left. \left. + \frac{98791\zeta_2}{972} + \frac{342\zeta_2\zeta_5}{5} + \frac{2978\zeta_2\zeta_3}{27} - \frac{40148\zeta_2^2}{405} + \frac{517\zeta_2^2\zeta_3}{5} + \frac{2167\zeta_2^3}{630} \right) \right] \\
& + C_F N_F \left[\epsilon^4 \left(-\frac{19296691}{7776} - 254\zeta_7 + \frac{22948\zeta_5}{45} + \frac{192068\zeta_3}{81} - 460\zeta_3^2 + \frac{75305\zeta_2}{648} \right. \right. \\
& \quad \left. \left. - 32\zeta_2\zeta_5 - \frac{5716\zeta_2\zeta_3}{27} + \frac{585929\zeta_2^2}{1620} - \frac{6724\zeta_2^2\zeta_3}{45} - \frac{2024\zeta_2^3}{105} \right) \right] \tag{10}
\end{aligned}$$

The unrenormalised three-loop quark form factor \mathcal{F}_3^q through to (and including) $\mathcal{O}(\epsilon^0)$ is given in eq. (5.4) of Ref. [16]. The pole contributions of \mathcal{F}_3^q are also given in eq. (3.7) of ref. [9] while the finite parts of the N_F^2 , $C_A N_F$ and $C_F N_F$ contributions are given in eq. (6) of ref. [10]. The finite $N_{F,V}$ contribution could already be inferred from [38]. The remaining finite contributions are also given in eqs. (8) and (9) of ref. [24]. The $\mathcal{O}(\epsilon^1)$ and $\mathcal{O}(\epsilon^2)$ contributions are given by,

$$\begin{aligned}
\mathcal{F}_3^q &= \mathcal{F}_3^q|_{\frac{1}{\epsilon^6}} + \dots + \mathcal{F}_3^q|_{\epsilon^0} \\
&+ C_F^3 \left[+ \epsilon \left(-\frac{343393}{48} - \frac{11896\zeta_7}{7} + \frac{22349\zeta_5}{3} + \frac{40835\zeta_3}{6} - 1203\zeta_3^2 - \frac{105553\zeta_2}{24} \right. \right. \\
&\quad \left. \left. - \frac{7858\zeta_2\zeta_5}{15} + \frac{6083\zeta_2\zeta_3}{6} + \frac{36693\zeta_2^2}{40} - \frac{3931\zeta_2^2\zeta_3}{6} + \frac{321227\zeta_2^3}{840} \right) \right. \\
&\quad + \epsilon^2 \left(-\frac{2512115}{96} + \frac{4160\zeta_{5,3}}{3} + \frac{45168\zeta_7}{7} + \frac{716537\zeta_5}{15} - \frac{137417\zeta_3}{12} \right. \\
&\quad \left. - \frac{33148\zeta_3\zeta_5}{3} + \frac{12749\zeta_3^2}{6} - \frac{797995\zeta_2}{48} - \frac{12361\zeta_2\zeta_5}{5} + \frac{18469\zeta_2\zeta_3}{2} \right. \\
&\quad \left. + 1985\zeta_2\zeta_3^2 + \frac{7653\zeta_2^2}{80} - \frac{15491\zeta_2^2\zeta_3}{20} + \frac{1147979\zeta_2^3}{240} - \frac{74208727\zeta_2^4}{50400} \right) \Big] \\
&+ C_F^2 C_A \left[+ \epsilon \left(+\frac{783459131}{34992} - 1349\zeta_7 - \frac{1894909\zeta_5}{270} - \frac{1259477\zeta_3}{54} + \frac{85649\zeta_3^2}{18} \right. \right. \\
&\quad + \frac{19394303\zeta_2}{1944} + \frac{4851\zeta_2\zeta_5}{5} - \frac{195175\zeta_2\zeta_3}{108} - \frac{15062939\zeta_2^2}{6480} \\
&\quad \left. + \frac{9751\zeta_2^2\zeta_3}{20} - \frac{1811231\zeta_2^3}{15120} \right) \\
&\quad + \epsilon^2 \left(+\frac{16308475427}{209952} - \frac{15472\zeta_{5,3}}{15} + \frac{415489\zeta_7}{42} - \frac{7913725\zeta_5}{162} - \frac{27356135\zeta_3}{324} \right. \\
&\quad + \frac{72904\zeta_3\zeta_5}{15} + \frac{2174933\zeta_3^2}{108} + \frac{521534243\zeta_2}{11664} + \frac{53128\zeta_2\zeta_5}{15} - \frac{5620115\zeta_2\zeta_3}{324} \\
&\quad - 1425\zeta_2\zeta_3^2 - \frac{161423233\zeta_2^2}{19440} + \frac{1083953\zeta_2^2\zeta_3}{180} - \frac{211343621\zeta_2^3}{90720} \\
&\quad \left. - \frac{22796551\zeta_2^4}{63000} \right) \Big] \\
&+ C_F C_A^2 \left[+ \epsilon \left(-\frac{458292965}{26244} - \frac{211\zeta_7}{18} + \frac{15601\zeta_5}{5} + \frac{42813461\zeta_3}{2916} - \frac{71734\zeta_3^2}{27} \right. \right. \\
&\quad - \frac{52068575\zeta_2}{8748} - \frac{1568\zeta_2\zeta_5}{9} + \frac{13139\zeta_2\zeta_3}{27} + \frac{4467743\zeta_2^2}{3240} - \frac{4408\zeta_2^2\zeta_3}{45} \\
&\quad \left. - \frac{8009\zeta_2^3}{945} \right) \\
&\quad + \epsilon^2 \left(-\frac{34868838031}{472392} - \frac{3592\zeta_{5,3}}{45} - \frac{176495\zeta_7}{36} + \frac{18727307\zeta_5}{810} + \frac{405838949\zeta_3}{5832} \right. \\
&\quad + \frac{568\zeta_3\zeta_5}{3} - \frac{820579\zeta_3^2}{54} - \frac{1546106255\zeta_2}{52488} - \frac{23456\zeta_2\zeta_5}{15} + \frac{2116327\zeta_2\zeta_3}{324} \\
&\quad \left. + \frac{2896\zeta_2\zeta_3^2}{9} + \frac{167549\zeta_2^2}{27} - 3805\zeta_2^2\zeta_3 + \frac{201469\zeta_2^3}{216} + \frac{6341548\zeta_2^4}{23625} \right) \Big]
\end{aligned}$$

$$\begin{aligned}
& +C_F^2 N_F \left[+ \epsilon \left(-\frac{50187205}{17496} + \frac{5863\zeta_5}{135} + \frac{929587\zeta_3}{243} - \frac{5771\zeta_3^2}{9} - \frac{1263505\zeta_2}{972} \right. \right. \\
& \quad \left. \left. - \frac{8515\zeta_2\zeta_3}{54} + \frac{821749\zeta_2^2}{3240} - \frac{875381\zeta_2^3}{7560} \right) \right. \\
& \quad + \epsilon^2 \left(-\frac{861740653}{104976} - \frac{294430\zeta_7}{63} + \frac{167299\zeta_5}{81} + \frac{32307433\zeta_3}{1458} - \frac{208487\zeta_3^2}{54} \right. \\
& \quad \left. - \frac{32868205\zeta_2}{5832} + \frac{953\zeta_2\zeta_5}{15} - \frac{152867\zeta_2\zeta_3}{162} + \frac{17061119\zeta_2^2}{9720} - \frac{172799\zeta_2^2\zeta_3}{180} \right. \\
& \quad \left. \left. - \frac{4769039\zeta_2^3}{6480} \right) \right] \\
& +C_F C_A N_F \left[+ \epsilon \left(+\frac{24570881}{4374} - \frac{28156\zeta_5}{45} - \frac{2418896\zeta_3}{729} + \frac{10816\zeta_3^2}{27} + \frac{7137385\zeta_2}{4374} \right. \right. \\
& \quad \left. \left. + \frac{2674\zeta_2\zeta_3}{27} - \frac{352559\zeta_2^2}{1620} + \frac{17324\zeta_2^3}{945} \right) \right. \\
& \quad + \epsilon^2 \left(+\frac{5509319623}{236196} + 1170\zeta_7 - \frac{622178\zeta_5}{135} - \frac{79031137\zeta_3}{4374} + \frac{218296\zeta_3^2}{81} \right. \\
& \quad \left. + \frac{102669593\zeta_2}{13122} + \frac{3272\zeta_2\zeta_5}{15} + \frac{11939\zeta_2\zeta_3}{81} - \frac{3829919\zeta_2^2}{3240} + \frac{9572\zeta_2^2\zeta_3}{15} \right. \\
& \quad \left. \left. + \frac{74461\zeta_2^3}{5670} \right) \right] \\
& +C_F N_F^2 \left[+ \epsilon \left(-\frac{2913928}{6561} + \frac{2248\zeta_5}{135} + \frac{2108\zeta_3}{27} - \frac{24950\zeta_2}{243} + \frac{68\zeta_2\zeta_3}{9} - \frac{3901\zeta_2^2}{810} \right) \right. \\
& \quad + \epsilon^2 \left(-\frac{109448624}{59049} + \frac{52828\zeta_5}{405} + \frac{848300\zeta_3}{2187} - \frac{1156\zeta_3^2}{81} - \frac{338858\zeta_2}{729} \right. \\
& \quad \left. \left. + \frac{1598\zeta_2\zeta_3}{27} - \frac{2573\zeta_2^2}{90} + \frac{44651\zeta_2^3}{5670} \right) \right] \\
& +C_F N_{F,V} \left(\frac{N^2 - 4}{N} \right) \times \\
& \left[+ \epsilon \left(+\frac{170}{3} + \frac{752\zeta_5}{9} + \frac{94\zeta_3}{9} - \frac{344\zeta_3^2}{3} + \frac{260\zeta_2}{3} + 30\zeta_2\zeta_3 - \frac{196\zeta_2^2}{15} - \frac{9728\zeta_2^3}{315} \right) \right. \\
& \quad + \epsilon^2 \left(+\frac{1460}{3} - \frac{4271\zeta_7}{3} + \frac{12970\zeta_5}{27} + \frac{2501\zeta_3}{27} - \frac{748\zeta_3^2}{9} + \frac{4345\zeta_2}{9} \right. \\
& \quad \left. \left. - \frac{256\zeta_2\zeta_5}{3} + \frac{239\zeta_2\zeta_3}{3} - \frac{3677\zeta_2^2}{45} - \frac{392\zeta_2^2\zeta_3}{3} + \frac{85244\zeta_2^3}{945} \right) \right] \tag{11}
\end{aligned}$$

Note that last colour factor is generated by graphs where the virtual gauge boson does not couple directly to the final-state quarks. This contribution is denoted by $N_{F,V}$ and is proportional to the charge weighted sum of the quark flavours. In the case of purely

electromagnetic interactions, we find,

$$N_{F,\gamma} = \frac{\sum_q e_q}{e_q}. \quad (12)$$

The unrenormalised three-loop gluon form factor through to (and including) $\mathcal{O}(\epsilon^0)$ is given in eq. (5.5) of Ref. [16]. The divergent parts are also given in eq. (8) of ref. [10] while the finite contributions are given in eq. (10) of ref. [24]. The $\mathcal{O}(\epsilon^1)$ and $\mathcal{O}(\epsilon^2)$ contributions for \mathcal{F}_3^g are given by,

$$\begin{aligned} \mathcal{F}_3^g &= \mathcal{F}_3^g|_{\frac{1}{\epsilon^6}} + \dots + \mathcal{F}_3^g|_{\epsilon^0} \\ &+ C_A^3 \left[+ \epsilon \left(+ \frac{270573319}{26244} - \frac{385579\zeta_7}{126} + \frac{389159\zeta_5}{135} - \frac{3601570\zeta_3}{729} + \frac{74899\zeta_3^2}{54} \right. \right. \\ &\quad \left. \left. - \frac{446863\zeta_2}{4374} + \frac{2449\zeta_2\zeta_5}{9} - \frac{34093\zeta_2\zeta_3}{54} - \frac{40819\zeta_2^2}{180} - \frac{47803\zeta_2^2\zeta_3}{180} \right. \right. \\ &\quad \left. \left. + \frac{7200127\zeta_2^3}{15120} \right) \right. \\ &\quad + \epsilon^2 \left(+ \frac{30151577675}{472392} + \frac{12392\zeta_{5,3}}{45} + \frac{2169431\zeta_7}{126} + \frac{3101341\zeta_5}{405} - \frac{59902487\zeta_3}{1458} \right. \\ &\quad \left. - \frac{89996\zeta_3\zeta_5}{15} + \frac{16453\zeta_3^2}{2} - \frac{108299125\zeta_2}{26244} - \frac{6897\zeta_2\zeta_5}{10} - \frac{80255\zeta_2\zeta_3}{27} \right. \\ &\quad \left. + \frac{7936\zeta_2\zeta_3^2}{9} - \frac{34875497\zeta_2^2}{9720} + \frac{714109\zeta_2^2\zeta_3}{360} + \frac{12226469\zeta_2^3}{5040} - \frac{1183759981\zeta_2^4}{756000} \right) \Big] \\ &+ C_A^2 N_F \left[+ \epsilon \left(- \frac{48658741}{8748} - \frac{10066\zeta_5}{45} + \frac{349918\zeta_3}{729} - \frac{11657\zeta_3^2}{27} + \frac{904045\zeta_2}{4374} \right. \right. \\ &\quad \left. \left. + \frac{791\zeta_2\zeta_3}{9} - \frac{34931\zeta_2^2}{1620} - \frac{52283\zeta_2^3}{1080} \right) \right. \\ &\quad + \epsilon^2 \left(- \frac{15039308929}{472392} - \frac{14271\zeta_7}{7} - \frac{391564\zeta_5}{405} + \frac{13422322\zeta_3}{2187} - \frac{76349\zeta_3^2}{81} \right. \\ &\quad \left. + \frac{66386911\zeta_2}{26244} + \frac{307\zeta_2\zeta_5}{5} + \frac{31849\zeta_2\zeta_3}{81} + \frac{373234\zeta_2^2}{1215} - \frac{104327\zeta_2^2\zeta_3}{180} \right. \\ &\quad \left. - \frac{6878021\zeta_2^3}{22680} \right) \Big] \\ &+ C_A C_F N_F \left[+ \epsilon \left(- \frac{10508593}{2916} + \frac{17092\zeta_5}{27} + \frac{240934\zeta_3}{243} + \frac{4064\zeta_3^2}{9} + \frac{8869\zeta_2}{54} \right. \right. \\ &\quad \left. \left. + \frac{640\zeta_2\zeta_3}{9} + \frac{28823\zeta_2^2}{270} + \frac{23624\zeta_2^3}{315} \right) \right. \\ &\quad + \epsilon^2 \left(- \frac{418631245}{17496} + \frac{16658\zeta_7}{9} + \frac{386102\zeta_5}{81} + \frac{4492979\zeta_3}{729} + \frac{17176\zeta_3^2}{27} \right. \\ &\quad \left. + \frac{163523\zeta_2}{108} - 496\zeta_2\zeta_5 + \frac{3500\zeta_2\zeta_3}{9} + \frac{437599\zeta_2^2}{540} + \frac{3148\zeta_2^2\zeta_3}{5} \right) \Big] \end{aligned}$$

$$\begin{aligned}
& + \frac{157424\zeta_2^3}{315} \Big) \Big] \\
& + C_F^2 N_F \left[+ \epsilon \left(+ \frac{18613}{54} - \frac{3080\zeta_5}{3} + \frac{10552\zeta_3}{9} - 272\zeta_3^2 - \frac{74\zeta_2}{3} - 16\zeta_2\zeta_3 + \frac{328\zeta_2^2}{5} - \frac{35648\zeta_2^3}{315} \right) \right. \\
& \quad + \epsilon^2 \left(+ \frac{383765}{162} - \frac{8828\zeta_7}{3} - \frac{35956\zeta_5}{9} + \frac{229772\zeta_3}{27} - \frac{6400\zeta_3^2}{3} - \frac{4109\zeta_2}{18} \right. \\
& \quad \left. \left. + 560\zeta_2\zeta_5 - 276\zeta_2\zeta_3 + 764\zeta_2^2 - \frac{1232\zeta_2^2\zeta_3}{3} - \frac{796168\zeta_2^3}{945} \right) \right] \\
& + C_A N_F^2 \left[+ \epsilon \left(+ \frac{16823771}{26244} + \frac{9368\zeta_5}{135} + \frac{5440\zeta_3}{27} - \frac{30283\zeta_2}{1458} - \frac{988\zeta_2\zeta_3}{27} + \frac{14018\zeta_2^2}{405} \right) \right. \\
& \quad + \epsilon^2 \left(+ \frac{1534229129}{472392} + \frac{33136\zeta_5}{81} + \frac{1698929\zeta_3}{2187} - \frac{17908\zeta_3^2}{81} - \frac{1822421\zeta_2}{8748} \right. \\
& \quad \left. \left. - \frac{15928\zeta_2\zeta_3}{81} + \frac{20009\zeta_2^2}{135} + \frac{12851\zeta_2^3}{5670} \right) \right] \\
& + C_F N_F^2 \left[+ \epsilon \left(+ \frac{196900}{243} - \frac{800\zeta_5}{9} - \frac{4208\zeta_3}{9} - 54\zeta_2 + \frac{112\zeta_2\zeta_3}{3} - \frac{2464\zeta_2^2}{45} \right) \right. \\
& \quad + \epsilon^2 \left(+ \frac{6322579}{1458} - \frac{17600\zeta_5}{27} - \frac{223756\zeta_3}{81} + \frac{3232\zeta_3^2}{9} - \frac{9626\zeta_2}{27} + \frac{2464\zeta_2\zeta_3}{9} \right. \\
& \quad \left. \left. - \frac{4913\zeta_2^2}{15} + \frac{248\zeta_2^3}{63} \right) \right] \tag{13}
\end{aligned}$$

The renormalised form factors are directly related to the unrenormalised form factors and details on how to extract the renormalised form factors to this order are given in section 2 of Ref. [16].

In this letter, we computed the three-loop quark and gluon form factors through to $\mathcal{O}(\epsilon^2)$ in the dimensional regularisation parameter. These contributions are relevant in the study of the infrared singularity structure at four loops. In particular, the $\mathcal{O}(\epsilon)$ terms of the three-loop form factors are required for the extraction of the four-loop quark and gluon collinear anomalous dimensions. The $\mathcal{O}(\epsilon^2)$ terms contribute to the finite part of the infrared-subtraction of the form factors at four loops. It is this infrared-subtracted finite part which is relevant for the study of the next-to-next-to-next-to-next-to-leading (N^4 LO) Drell-Yan and Higgs production processes. In particular, the $\mathcal{O}(\epsilon^2)$ three-loop contributions represent a finite ingredient to these processes at four-loops.

Acknowledgements

We thank Daniel Maître and Volodya Smirnov for useful comments regarding the conversion of $\zeta_{-6,-2}$ to $\zeta_{5,3}$. This research was supported in part by the Swiss National Science Foundation (SNF) under contract 200020-126691, by the Forschungskredit der Universität Zürich, the UK Science and Technology Facilities Council, by the European Commission's

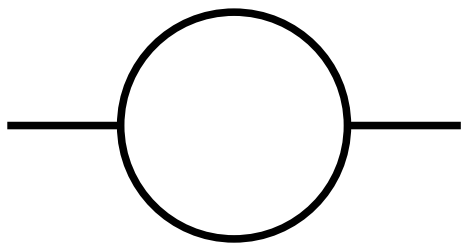
Marie-Curie Research Training Network under contract MRTN-CT-2006-035505 ‘Tools and Precision Calculations for Physics Discoveries at Colliders’, by the Helmholtz Alliance “Physics at the Terascale”, and by Deutsche Forschungsgemeinschaft (DFG SCHR 993/2-1). EWNG gratefully acknowledges the support of the Wolfson Foundation and the Royal Society.

References

- [1] G. Altarelli, R. K. Ellis and G. Martinelli, Nucl. Phys. B **157** (1979) 461.
- [2] R. Hamberg, W. L. van Neerven and T. Matsuura, Nucl. Phys. B **359** (1991) 343; B **644** (2002) 403(E).
- [3] R.V. Harlander and W. B. Kilgore, Phys. Rev. Lett. **88** (2002) 201801 [hep-ph/0201206].
- [4] S. Dawson, Nucl. Phys. B **359** (1991) 283.
- [5] D. Graudenz, M. Spira and P.M. Zerwas, Phys. Rev. Lett. **70** (1993) 1372;
M. Spira, A. Djouadi, D. Graudenz and P. M. Zerwas, Nucl. Phys. B **453** (1995) 17 [hep-ph/9504378];
A. Djouadi, M. Spira and P.M. Zerwas, Z. Phys. C **70** (1996) 427 [hep-ph/9511344];
M. Spira, Fortsch. Phys. **46** (1998) 203 [hep-ph/9705337];
R. Harlander and P. Kant, JHEP **0512** (2005) 015 [hep-ph/0509189].
- [6] S. Catani, D. de Florian and M. Grazzini, JHEP **0105** (2001) 025 [hep-ph/0102227];
C. Anastasiou and K. Melnikov, Nucl. Phys. B **646** (2002) 220 [hep-ph/0207004];
V. Ravindran, J. Smith and W. L. van Neerven, Nucl. Phys. B **665** (2003) 325 [hep-ph/0302135], **704** (2005) 332 [hep-ph/0408315].
- [7] S. Moch, J. A. M. Vermaseren and A. Vogt, Nucl. Phys. B **688** (2004) 101 [hep-ph/0403192].
- [8] A. Vogt, S. Moch and J. A. M. Vermaseren, Nucl. Phys. B **691** (2004) 129 [hep-ph/0404111].
- [9] S. Moch, J. A. M. Vermaseren and A. Vogt, JHEP **0508** (2005) 049 [hep-ph/0507039].
- [10] S. Moch, J. A. M. Vermaseren and A. Vogt, Phys. Lett. B **625**, 245 (2005) [hep-ph/0508055].
- [11] G. Kramer and B. Lampe, Z. Phys. C **34** (1987) 497; **42** (1989) 504(E);
T. Matsuura and W.L. van Neerven, Z. Phys. C **38** (1988) 623;
T. Matsuura, S.C. van der Maarck and W.L. van Neerven, Nucl. Phys. B **319** (1989) 570.
- [12] R.V. Harlander, Phys. Lett. B **492** (2000) 74 [hep-ph/0007289].

- [13] T. Gehrmann, T. Huber and D. Maître, Phys. Lett. B **622** (2005) 295 [hep-ph/0507061].
- [14] T. Becher, M. Neubert and B. D. Pecjak, JHEP **0701** (2007) 076 [hep-ph/0607228].
- [15] T. Becher and M. Neubert, JHEP **0906** (2009) 081 [0903.1126].
- [16] T. Gehrmann, E. W. N. Glover, T. Huber, N. Ikizlerli and C. Studerus, JHEP **1006** (2010) 094 [1004.3653].
- [17] K.G. Chetyrkin and F.V. Tkachov, Nucl. Phys. **B192** (1981) 159.
- [18] S.G. Gorishnii, S.A. Larin, L.R. Surguladze and F.V. Tkachov, Comput. Phys. Comm. **55** (1989) 381;
S.A. Larin, F.V. Tkachov and J.A.M. Vermaseren, NIKHEF-H-91-18.
- [19] S. Bekavac, Comput. Phys. Commun. **175** (2006) 180 [hep-ph/0505174].
- [20] T. Gehrmann, G. Heinrich, T. Huber and C. Studerus, Phys. Lett. B **640** (2006) 252 [hep-ph/0607185].
- [21] G. Heinrich, T. Huber and D. Maître, Phys. Lett. B **662** (2008) 344 [0711.3590].
- [22] G. Heinrich, T. Huber, D. A. Kosower and V. A. Smirnov, Phys. Lett. B **678** (2009) 359 [0902.3512].
- [23] R.N. Lee, A.V. Smirnov and V.A. Smirnov, JHEP **1004** (2010) 020 [1001.2887].
- [24] P. A. Baikov, K. G. Chetyrkin, A. V. Smirnov, V. A. Smirnov and M. Steinhauser, Phys. Rev. Lett. **102** (2009) 212002 [0902.3519].
- [25] R. N. Lee and V. A. Smirnov, [1010.1334].
- [26] O. V. Tarasov, Phys. Rev. D **54** (1996) 6479 [hep-th/9606018].
- [27] R. N. Lee, A. V. Smirnov and V. A. Smirnov, [1005.0362].
- [28] T. Huber and D. Maître, Comput. Phys. Commun. **178** (2008) 755 [0708.2443].
- [29] M. Czakon, Comput. Phys. Commun. **175** (2006) 559 [hep-ph/0511200].
- [30] A. V. Smirnov and M. N. Tentyukov, Comput. Phys. Commun. **180** (2009) 735 [0807.4129].
- [31] A. V. Smirnov, V. A. Smirnov and M. Tentyukov, [0912.0158].
- [32] T. Huber and D. Maître, Comput. Phys. Commun. **175** (2006) 122 [hep-ph/0507094].
- [33] J. Blümlein, D. J. Broadhurst and J. A. M. Vermaseren, Comput. Phys. Commun. **181** (2010) 582 [0907.2557].
- [34] D. Maître, Comput. Phys. Commun. **174** (2006) 222 [hep-ph/0507152].

- [35] F. Wilczek, Phys. Rev. Lett. **39** (1977) 1304;
M.A. Shifman, A.I. Vainshtein and V.I. Zakharov, Phys. Lett. B **78** (1978) 443;
J.R. Ellis, M.K. Gaillard, D.V. Nanopoulos and C.T. Sachrajda, Phys. Lett. B **83**
(1979) 339;
T. Inami, T. Kubota and Y. Okada, Z. Phys. C **18** (1983) 69.
- [36] K.G. Chetyrkin, B.A. Kniehl and M. Steinhauser, Phys. Rev. Lett. **79** (1997) 353
[hep-ph/9705240].
- [37] B.A. Kniehl and M. Spira, Z. Phys. C **69** (1995) 77 [hep-ph/9505225];
K.G. Chetyrkin, B.A. Kniehl and M. Steinhauser, Nucl. Phys. B **510** (1998) 61 [hep-
ph/9708255].
- [38] J. A. M. Vermaseren, A. Vogt and S. Moch, Nucl. Phys. B **724** (2005) 3 [hep-
ph/0504242].



b21

